### RESEARCH PAPER

# Red shift in absorption edge of $Cd_{1-x}Ni_xS$ dilute magnetic semiconductor nanofilms

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**Abstract** Nickel-doped cadmium sulphide dilute semiconductor nanofilms  $(Cd_{1-x}Ni_xS, 0 \le x \le 0.09)$  have been investigated using atomic force microscopy, scanning electron microscopy, Fourier infrared transform and UV–Vis spectroscopy. With increasing Ni content a decrease in average crystallite size (26-10 nm) and surface roughness has been found. A red shift in optical band gap has been observed with decrease in average crystallite size, which is against quantum confinement.

 $\begin{tabular}{ll} \textbf{Keywords} & \end{tabular} & \end{tabul$ 

#### Introduction

In the present research scenario of devising multifunctional materials, dilute magnetic semiconductors (DMS) have acquired an indispensable place in the field of spintronics, non-volatile memories, quantum computing, communication in the solid state, magneto-optical devices etc. (Dietl et al. 2000; Dietl 2008; 2012; Chandramohan et al. 2009; Kim et al. 1993). In this communication, we report the investigation of  $Cd_{1-x}Ni_xS$  ( $0 \le x \le 0.09$ ) DMS nanofilms via FTIR, SEM, AFM and UV–Vis spectroscopy. We observed and studied for the structure, surface morphology and optical band gap.

Ohno 2010; Radovanovic et al. 2005; Chelikowsky

et al. 2006; Jungwirth et al. 2006). DMS are produced

by replacing some of the cations in the host semiconductor (especially group II-VI and III-V) with

magnetic ions particularly transition metal (TM) like

Fe, Mn, Ni, Co etc. (Furdyna 1988; Malguth et al.

2008; Wu et al. 2006; Schwartz et al. 2004; Ozer et al.

2012). The spin-exchange (sp-d) interaction between

the semiconductor charge carriers and magnetic ions

add local magnetic moments in DMS and accumulate magneto-opto-electronic character in them. The intro-

duction of dopant in II-VI semiconductors is an

extensive route to tune them with respect to the phase, morphology and crystallite size. TM-doped CdS films provide an opportunity to integrate various properties into a single material for advanced applications (Chelikowsky et al. 2006; Wu et al. 2006; Saravanan

et al. 2011; Kamruzzaman et al. 2012; Polat et al.

Experimental

A detailed experimental procedure for deposition of  $Cd_{1-x}Ni_xS$  ( $0 \le x \le 0.09$ ) DMS nanofilms is given elsewhere (Kumar et al. 2013; Kumar et al. 2012a,

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2012b). The thickness of nanofilms (Table 1) was measured by stylus profilometer (AMBIOS XP-1). Surface morphology was studied using field emission scanning electron microscope (FE-SEM, HITACHI S-4700) and AFM (NTMDT-NTEGRA) in semicontact mode. FTIR (Perkin Elmer 1600 FTIR spectrophotometer) measurements were obtained in the spectral range 400–4000 cm<sup>-1</sup>. UV–Vis–NIR double beam spectrophotometer (Perkin Elmer, Lambda-750) was used to obtained transmittance (%T) and reflectance (%R) spectra in the wavelength range 300–900 nm.

#### Results and discussion

The structural analysis of  $Cd_{1-x}Ni_xS$  DMS nanofilms has been performed using XRD and given elsewhere (Kumar et al. 2013). All the nanofilms have prominent reflection peaks associated with  $\alpha$ -CdS structure in conjunction with other low-intensity reflection peaks assigned either  $\alpha$ -CdS or both  $\alpha$  and  $\beta$ -CdS (Kumar et al. 2013).

FTIR spectra for  $Cd_{1-x}Ni_xS$  nanofilms (Fig. 1) show various absorption bands/peaks at different frequencies. The absorption bands/peaks at  $\sim 3400-3430~cm^{-1}$ ,  $1438-1450~cm^{-1}$ ,  $860-900~cm^{-1}$ ,  $1370-1400~cm^{-1}$  and  $1990-2150~cm^{-1}$  and weak doublet at 2920 and  $2851~cm^{-1}$  have been assigned for O-H stretching, asymmetric scissor deformation ( $\delta_{as}$ -CH<sub>2</sub>), hydrogen bound O-H out of plane bending, primary or secondary OH in-plane bending, isothiocyanate (-NCS) formed due to thiourea dissociation,

asymmetrical and symmetrical vibration of CH<sub>2</sub> group ( $v_{as}$ -CH<sub>2</sub> and  $v_{s}$ -CH<sub>2</sub>), respectively (Coates 2000; Thangadurai et al. 2008; Cabana et al. 2011). The peak at ~1011 cm<sup>-1</sup> (x=0) and ~1075 cm<sup>-1</sup> (x=0.02–0.09) belongs to primary amine C–N stretch (Coates 2000) and also shared by C–O stretching (Thangadurai et al. 2008). However, secondary amine C–N stretch is indicated by peaks at 1399 cm<sup>-1</sup> (x=0) and ~1370 cm<sup>-1</sup> (x=0.02–0.09) (Coates 2000). These peaks of

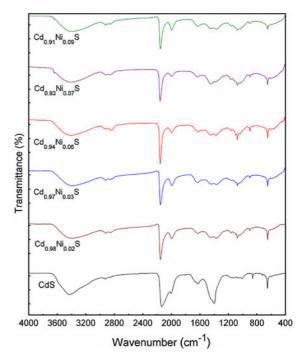


Fig. 1 FTIR spectra for  $Cd_{1-x}Ni_xS$  (0  $\leq x \leq$  0.09) nanofilms (The ordinate scale for different x values is shifted for clarity)

**Table 1** The compositional, morphological and optical parameters of  $Cd_{1-x}Ni_xS$  ( $0 \le x \le 0.09$ ) nanofilms, where  $x_m = \left(\frac{Ni}{Cd+Ni}\right)_{Solution}$  is the solution concentration from molarities and  $x_f = \left(\frac{Ni}{Cd+Ni}\right)_{Film}$  is the film composition from EDAX

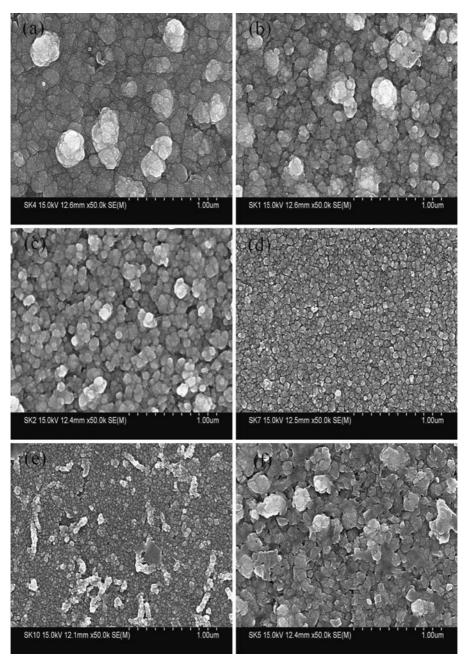
Sample	$x_{ m m}$	$x_{ m f}$	Film thickness (nm)	AFM		$\alpha (\times 10^4 \text{ cm}^{-1})$	$E_{\rm u}$
				D <sub>AFM</sub> (nm)	R <sub>f</sub> (nm)		(eV)
CdS	0	0	67.6	26	9.6	1.01	0.282
Cd <sub>0.98</sub> Ni <sub>0.02</sub> S	0.05	0.02	79.4	21	5.2	2.65	0.291
Cd <sub>0.97</sub> Ni <sub>0.03</sub> S	0.09	0.03	83.5	17	5.1	2.73	0.332
Cd <sub>0.94</sub> Ni <sub>0.06</sub> S	0.13	0.06	88.4	15	5.0	2.77	0.372
Cd <sub>0.93</sub> Ni <sub>0.07</sub> S	0.17	0.07	92.2	15	3.1	5.32	0.490
$Cd_{0.91}\ Ni_{0.09}S$	0.20	0.09	103.8	10	4.5	6.80	0.509



C–N stretching are evolved due to alkaline nature of solution because of the presence of TEA and NH<sub>3</sub>.  $Cd_{1-x}Ni_xS$  ( $x \ge 0.03$ ) nanofilms have small peaks in the region 470–600 cm<sup>-1</sup> belonging to S–S stretch (Coates 2000). All nanofilms have a strong absorption peak at  $\sim 650$  cm<sup>-1</sup> associated with Cd–S stretching (Thangadurai et al. 2008). No peak for Ni–S stretching has been observed which indicates homogeneous

substitution of Ni<sup>2+</sup> ions in CdS structure. The shift of Cd–S stretch ( $\sim$ 650 cm<sup>-1</sup>) towards lower wavenumber with increasing x indicates the incorporation of Ni<sup>2+</sup> ions in the CdS structure.

Figure 2 shows micro-structural surface morphology for  $Cd_{1-x}Ni_xS$  nanofilms at a scale of 1  $\mu$ m (50 K×). All nanofilms show uniformly distributed spherical nanocrystallites on the surface of substrate



**Fig. 2** SEM micrographs for  $Cd_{1-x}Ni_xS$  nanofilms (**a** x=0, **b** x=0.02, **c** x=0.03, **d** x=0.06, **e** x=0.07 and **f** x=0.09)

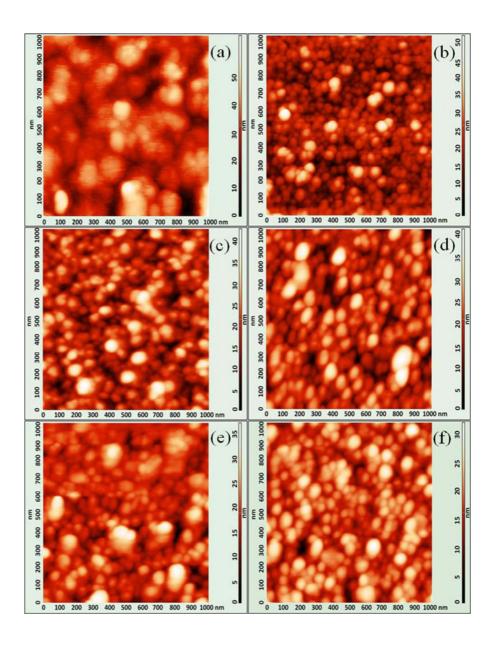


with good adhesion. The surface of all films is compact, densely packed, continuous and barren free which may be attributed to chemical activity of triton (TX-100) (Kumar et al. 2012a). Surface morphology of nanofilms is strongly dependent on Ni content. The substitution of Ni<sup>2+</sup> ions in CdS structure produces substantial surface changes with less agglomeration of the crystallites i.e. reduction in crystallite size and increase in grain boundaries. The nanofilm x = 0.07 shows sphere-like crystallite chains in an irregular spatial distribution. However, surface morphology for

x = 0.09 is quite different, showing sheet-like structural growth.

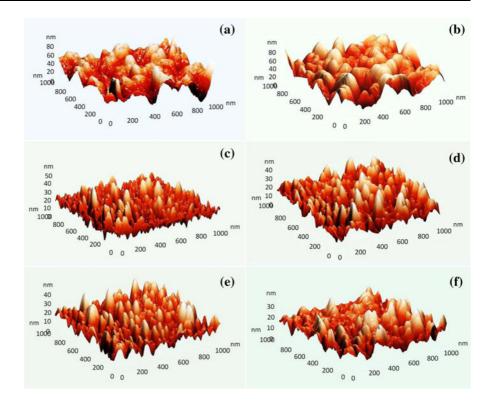
The surface topography and structure of film surfaces are shown in Fig. 3 (AFM-2D images) and Fig. 4 (AFM-3D images) of DMS nanofilms. Nanofilm with x = 0 (Fig. 3a) possesses larger crystallites  $(D_{AFM})$  and high average surface roughness  $(R_f)$  (Table 1). However, with increasing Ni<sup>2+</sup> concentration (Fig. 3a–f) the  $D_{AFM}$  of  $Cd_{1-x}Ni_xS$  nanofilms decreases.  $R_f$  also decreases up to  $x \le 0.07$ . The low value of surface skewness  $(S_{sk} < 0.7)$  and kurtosis

**Fig. 3** AFM-2D images for  $Cd_{1-x}Ni_xS$  nanofilms (**a** x = 0, **b** x = 0.02, **c** x = 0.03, **d** x = 0.06, **e** x = 0.07 and **f** x = 0.09)





**Fig. 4** AFM-3D images for  $Cd_{1-x}Ni_xS$  nanofilms (**a** x = 0, **b** x = 0.02, **c** x = 0.03, **d** x = 0.06, **e** x = 0.07 and **f** x = 0.09)



coefficient ( $S_{\rm ka} \geq 3$ ) for all nanofilms indicate that the height distribution is uniform, with approximately equal number of high peaks to deep valleys (Fig. 4a–f) over the scanned area (1  $\mu$ m  $\times$  1  $\mu$ m). All films grow with columnar structures along the c-axis perpendicular to the substrate which confirm the prominent hexagonal-wurtzite structure in deposited  $Cd_{1-x}Ni_xS$  nanofilms (Lazos et al. 2008).

Figure 5 shows %T and %R spectra of DMS nanofilms grown on glass substrate. The occurrence of low %T interference pattern in spectra indicates that the deposited nanofilms are ultra thin, smooth and homogeneous. High %T (>65 %) and low %R (<20 %) in the visible region makes these films good candidate for solar energy applications.

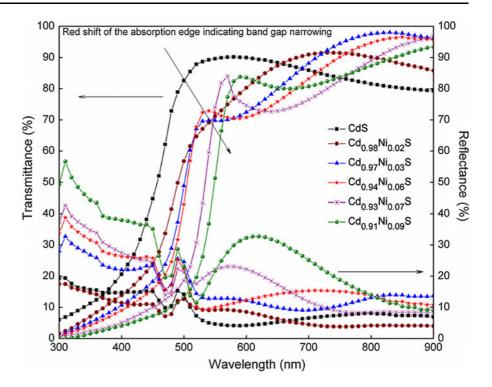
The optical behaviour of deposited DMS nanofilms have been analyzed in term of optical absorption coefficient ( $\alpha$ ) (Salem et al. 2008). The value of  $\alpha$  for all nanofilms has been found to be  $\sim 10^5$  cm<sup>-1</sup> in the visible region, and at absorption edge,  $\alpha$  increases with increasing x (Table 1). This may be due to decrease in crystallite size which provides large collective surface area and large number of absorption/scattering centres for light. The exponential dependence of  $\alpha$  on photon energy (hv) near the band edge has been studied in

terms of Urbach energy  $(E_{\rm u})$  by Urbach–Martienssen model (Pejova 2010).  $E_{\rm u}$  is often interpreted as the width of the tail of localized states in the gap region (for the values of  $E_{\rm u}$  see Table 1). A systematic increase in  $E_{\rm u}$  may be attributed to the creation of more localized states within the band tails of valence and conduction band due to the existence of defects and disorders. A sharp fall in %T (Fig. 5) near the fundamental absorption edge indicates a direct energy transition in the forbidden gap. The optical band gap  $(E_{\rm g})$  for direct transition has been determined using Tauc's relation (Kumar et al. 2012b). The  $E_{\rm g}$  values have been estimated by extrapolating  $(\alpha hv)^2 \rightarrow 0$  (Fig. 6). The band gap decreases with increasing Ni content in CdS nanofilms (Fig. 7).

The variation of  $E_{\rm g}(x_{\rm f})$  with film composition  $(x_{\rm f})$  (inset Fig. 7) indicates non-linear behaviour showing bowing phenomenon following Vegard's law (Li et al. 2011). The red shift in band edge on Ni doping (Fig. 6) indicates narrowing of band gap even with decrease in crystallite size. This may be attributed to structural disorders, increased tail width of localized states  $(E_{\rm u})$  (Table 1) and sp-d hybridization effect. Ni<sup>2+</sup> ions in the host CdS crystal form new localized electronic states in the band gap, which arise from the 3d-shell of



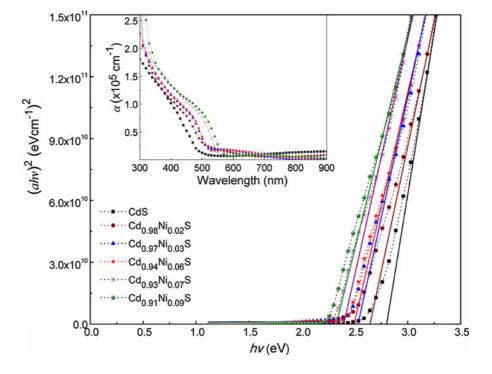
**Fig. 5** %*T* and %*R* spectra for  $Cd_{1-x}Ni_xS$  nanofilms  $(0 \le x \le 0.09)$ 



 $Ni^{2+}$  ion under the action of surrounding CdS crystal field (Podlowski et al. 1992). With an increase in  $Ni^{2+}$  ion concentration, the sp-d exchange interaction between the band electrons and the localized

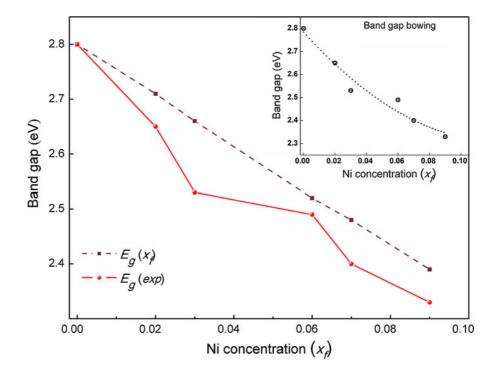
d-electrons of Ni<sup>2+</sup> ions increases (Polat et al. 2012; Kim et al. 1993). The strength of this interaction strongly depends on the number of d-electrons (Samanta et al. 2009). The reduction of crystallite

**Fig. 6** Tauc's plot for  $Cd_{1-x}Ni_xS$  ( $0 \le x \le 0.09$ ) nanofilms and inset show  $\alpha$  *vs.* wavelength





**Fig. 7** Variation of  $E_g$  with Ni concentration  $(x_f)$  and inset show band gap bowing



size with doping of Ni<sup>2+</sup> ions also leads to an increase in surface/volume ratio. As a result, the surface states corresponding to Ni<sup>2+</sup> in CdS increase and reduce the excitonic emission via non-radiative surface recombination (Podlowski et al. 1992). The Ni doping creates high density of impurity states in the nanofilms, which may also perturbate the band structure in the energy gap. The impurity band merges with the nearest intrinsic band and the Fermi level may lie inside the parabolic portion of the valence band (Oztas et al. 2007). Thus, less energy will be required for the electrons to move from the Fermi level into the conduction band (Oztas et al. 2007). The dopant ions form deep trap levels, which act as luminescence centres, create more discrete energy states by modifying the band structure.

## Conclusion

The morphology and optical behaviour of II–VI DMS have been reported.  $\mathrm{Ni}^{2+}$  ions in CdS show change in absorption peak position ( $\sim 650~\mathrm{cm}^{-1}$ ) and the strength of respective vibrational bands indicate the incorporation of Ni in host CdS lattice. The surface morphology depends on  $\mathrm{Ni}^{2+}$  concentration and

shows less agglomeration of the crystallites with increase in grain boundaries and decrease in surface roughness on an increase in Ni content. The red shift in band edge on incorporation of  $\mathrm{Ni}^{2+}$  ions in CdS show band gap narrowing with decrease in crystallite size. This may be attributed to creation of localized states within the band tails and  $\mathit{sp-d}$  hybridization effect. The high transmittance and band gap tailoring of DMS nanofilms from 2.80 to 2.33 eV in visible region and UV end of solar spectrum make them suitable for large area coating of photovoltaic devices, anti-reflection coating, solar control coating and warming window layers.

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